

UDC 517

SCOPUS CODE 2201

<https://doi.org/10.36073/1512-0996-2024-1-170-175>

Investigation of Pericyclic Mechanism's Kinematics

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Abstract. A kinematic analysis of the pericyclic mechanism, where the gear is internally engaged with fixed central wheel is presented. Geometric features the trajectory of the mechanism allows use of this type of mechanisms in machines, which provides such kind of technological processes where slight fluctuations of the driven link sometimes are necessary at the extreme position of the system. Mechanisms of this type are used in the light industry, where the processing of any tissue requires slight fluctuations of the executive body at the temporary position.

Keywords: cycloid; force; mechanism; movement; satellite; work.

Introduction

Planetary gears with a satellite or with satellites, which allow an approximate quasi-stop of the output link, are widely used in textile and light industry machines. Such allowances for minor deviations during the dwell of the working body of the mechanism are due to the flexibility and pliability of the materials processed and driven in this industry. For example, in special sewing machines, there are mechanisms for transverse or longitudinal vibrations of the needle when stitching the material. This behavior of the tool does not have a damaging effect on the material being processed, in the process of puncturing the material with a needle. In such cases, it is possible to use mechanisms with less kinematic complexity and making it

possible to develop as high speeds as possible for high-speed sewing and textile machines.

Main Part

In this paper, we consider a planetary gear with a planetary gear that is in internal engagement with a stationary central wheel (Fig. 1). Our task is to determine the speed and acceleration of the hole for the sewing thread of the driven link when the driving link $O_1 O_2$ rotates with a certain angular velocity and angular acceleration. The angular velocities of the carrier and the moving link are related by the relation:

$$\omega_1 \cdot |O_1 O_2| = \omega_2 \cdot |P O_2| \quad (1)$$

whence the angular velocity of the link 2 will be:

$$\omega_2 = \frac{|O_1 O_2|}{|P O_2|} \cdot \omega_1 = \frac{r_2 - r_1}{r_2} \cdot \omega_1 \quad (2)$$

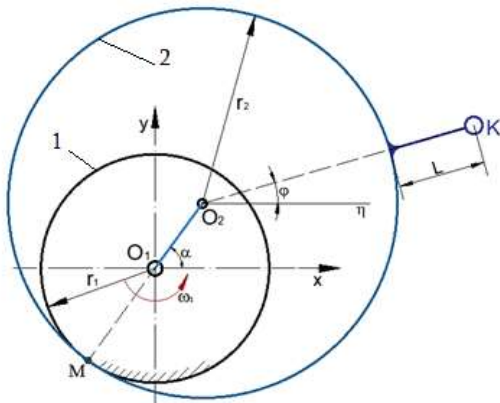


Fig. 1

To determine the speeds of any point of the driven link, it is necessary to find the position of the instantaneous center of speed of the gear with internal gearing. The mentioned, instantaneous center of speed is the point P of contact between gear 2 and wheel of gear 1.

The crank $O_1 O_2$ moves with an angular velocity ω_1 and with an angular acceleration ε_1 relative to the point O_1 . Based on this, the speed of the point O_2 is determined by the formula:

$$V_{O_2} = \omega_1 \cdot |O_1 O_2| = (r_2 - r_1) \cdot \omega_1. \quad (3)$$

After that, we find the speed of point K by the following formula:

$$V_{K O_2} = \omega_2 \cdot |K O_2| = \frac{(r_2 - r_1) \cdot \dot{\alpha}}{r_2} \cdot (r_2 + L). \quad (4)$$

Full velocity of the point K, can be written as:

$$\vec{V}_K = \vec{V}_{O_2} + \vec{V}_{K O_2}; \quad (5)$$

Projections of the full velocity of the point K on the axes x and y will be written as:

$$\begin{aligned} V_{Kx} &= V_{O_2} \cdot \cos(90^\circ + \alpha) + V_{K O_2} \cdot \cos(90^\circ + \varphi) = \\ &= -(r_2 - r_1) \cdot \dot{\alpha} \cdot \sin \alpha - \frac{(r_2 - r_1) \cdot \dot{\alpha}}{r_2} \cdot (r_2 + L) \cdot \\ &\quad \sin(1 - r_1/r_2) \alpha; \end{aligned} \quad (6)$$

$$\begin{aligned} V_{Ky} &= V_{O_2} \cdot \sin(90^\circ + \alpha) + V_{K O_2} \cdot \sin(90^\circ + \varphi) = \\ &= (r_2 - r_1) \cdot \dot{\alpha} \cdot \cos \alpha + \frac{(r_2 - r_1) \cdot \dot{\alpha}}{r_2} \cdot (r_2 + L) \cdot \\ &\quad \cos(1 - r_1/r_2) \alpha; \end{aligned} \quad (7)$$

And whence we will have:

$$V_K = \sqrt{V_{Kx}^2 + V_{Ky}^2}. \quad (8)$$

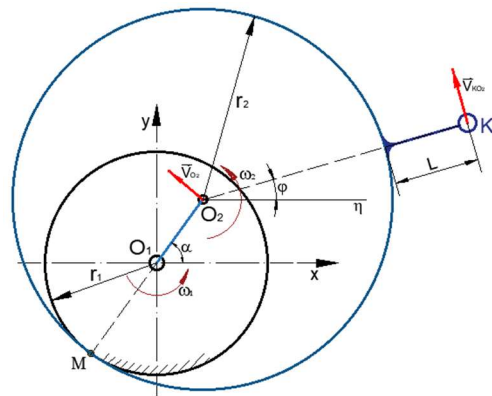


Fig. 2

At the same time, if the crank $O_1 O_2$ rotates around the point O_1 with a constant angular velocity ω_1 then the centripetal acceleration $W_{O_2}^n$ will have a view:

$$W_{O_2}^n = \omega_1^2 \cdot (r_2 - r_1) = \dot{\alpha}^2 \cdot (r_2 - r_1). \quad (8)$$

Centripetal acceleration of point K of the gear in relative motion can be determined as:

$$W_K^n = \omega_2^2 \cdot (r_2 + L) = \frac{(r_2 - r_1)^2 \cdot (r_2 + L)}{r_2^2} \cdot \dot{\alpha}^2. \quad (9)$$

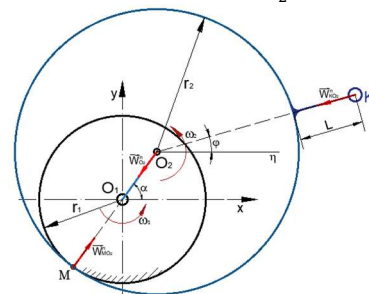


Fig. 3

The vector of full acceleration of point K is found as the sum of vectors of translational and relative acceleration.

$$\vec{W}_K = \vec{W}_{O_2}^n + \vec{W}_{KO_2}^n. \quad (10)$$

At the same time the meaning of modules of here given accelerations are equal.

$$\begin{aligned} W_{Kx} &= -W_{O_2}^n \cdot \cos\alpha - W_{KO_2}^n \cdot \cos\varphi = \\ &= -\dot{\alpha}^2 \cdot (r_2 - r_1) \cdot \cos\alpha - \frac{(r_2 - r_1)^2 \cdot (r_2 + L)}{r_2^2} \cdot \dot{\alpha}^2 \cdot \\ &\quad \cos(1 - r_1/r_2)\alpha; \end{aligned} \quad (11)$$

$$\begin{aligned} W_{Ky} &= -W_{O_2}^n \cdot \sin\alpha - W_{KO_2}^n \cdot \sin\varphi = \\ &= -\dot{\alpha}^2 \cdot (r_2 - r_1) \cdot \sin\alpha - \frac{(r_2 - r_1)^2 \cdot (r_2 + L)}{r_2^2} \cdot \dot{\alpha}^2 \cdot \\ &\quad \sin(1 - r_1/r_2)\alpha. \end{aligned} \quad (12)$$

The angle among them is:

$$W_K = \sqrt{W_{Kx}^2 + W_{Ky}^2}. \quad (13)$$

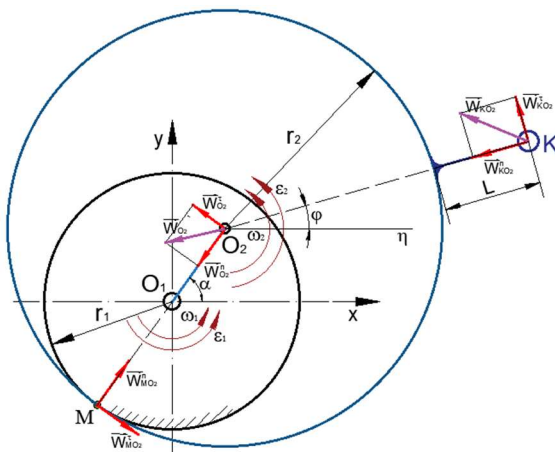


Fig. 4

If the crank, in addition to the angular velocity, also has an angular acceleration, then to determine the total acceleration of the point K, it is necessary to use the following analytical dependence.

$$\omega_2 \cdot r_2 = (r_2 - r_1)\dot{\alpha}; \quad (14)$$

$$\omega_2 = \frac{(r_2 - r_1)\dot{\alpha}}{r_2}. \quad (15)$$

Differentiating this fraction with respect of time, we find the angular acceleration of the gear:

$$\varepsilon_2 = \frac{d\omega_2}{dt} = \frac{(r_2 - r_1)\ddot{\alpha}}{r_2}. \quad (16)$$

The tangential acceleration of the point O2, as the point of the crank O1 O2, is directed towards O1 and is equal in absolute value to:

$$W_{O_2}^n = \dot{\alpha}^2 \cdot (r_2 - r_1). \quad (17)$$

The module of rotation acceleration of the point O2 is perpendicular to the straight-line O1O2 can be calculated as:

$$W_{O_2}^t = \ddot{\alpha} \cdot (r_2 - r_1). \quad (18)$$

If the point O2 is taken as the pole, then the total acceleration of the point K can be found by the following equation. Based on this we will have:

$$\vec{W}_K = \vec{W}_{O_2} + \vec{W}_{KO_2}^n + \vec{W}_{KO_2}^t. \quad (19)$$

The acceleration of the point K has already been found through the components $W_{O_2}^n$ and $W_{O_2}^t$. The tangent acceleration of the point K relative to the point O2 is directed towards O2 the module of this acceleration is equal to $\omega_2^2 \cdot (r_2 + L)$. Substituting the value ω_2 , we will have:

$$W_{KO_2}^n = \omega_2^2 \cdot (r_2 + L) = \frac{(r_2 - r_1)^2 \cdot (r_2 + L)}{r_2^2} \cdot \dot{\alpha}^2. \quad (20)$$

The rotational acceleration of the point K relative to the point O2 is directed perpendicular to the straight-line O2K and the algebraic value of this acceleration is:

$$W_{KO_2}^t = \varepsilon_2 \cdot (r_2 + L) = \frac{(r_2 - r_1) \cdot (r_2 + L) \cdot \ddot{\alpha}}{r_2}; \quad (21)$$

Projecting both sides of equality (22) onto the x and y axes

$$\begin{aligned} W_{Kx} &= -W_{O_2O_1}^n \cdot \cos\alpha - W_{O_2O_1}^t \cdot \sin\alpha - W_{KO_2}^n \cdot \\ &\quad \cos\varphi - W_{KO_2}^t \cdot \sin\varphi; \\ W_{Ky} &= -W_{O_2O_1}^n \cdot \sin\alpha + W_{O_2O_1}^t \cdot \cos\alpha - W_{KO_2}^n \cdot \\ &\quad \sin\varphi + W_{KO_2}^t \cdot \cos\varphi. \end{aligned} \quad (22)$$

In expanded form, these equations can be written as:

$$\begin{aligned} W_{Kx} &= -(r_2 - r_1) \cdot \dot{\alpha}^2 \cdot \cos\alpha - (r_2 - r_1) \cdot \ddot{\alpha} \sin\alpha - \\ &\quad - \frac{(r_2 - r_1)^2 \cdot (r_2 + L)}{r_2^2} \cdot \dot{\alpha}^2 \cdot \cos(1 - r_1/r_2)\alpha - \\ &\quad - \frac{(r_2 - r_1) \cdot (r_2 + L) \cdot \ddot{\alpha}}{r_2} \cdot \sin(1 - r_1/r_2)\alpha; \\ W_{Ky} &= -(r_2 - r_1) \cdot \dot{\alpha}^2 \cdot \sin\alpha + (r_2 - r_1) \cdot \ddot{\alpha} \cos\alpha - \\ &\quad - \frac{(r_2 - r_1)^2 \cdot (r_2 + L)}{r_2^2} \cdot \dot{\alpha}^2 \cdot \sin(1 - r_1/r_2)\alpha + \\ &\quad + \frac{(r_2 - r_1) \cdot (r_2 + L) \cdot \ddot{\alpha}}{r_2} \cdot \cos(1 - r_1/r_2)\alpha. \end{aligned} \quad (23)$$

The full acceleration value is calculated as follows:

$$W_K = \sqrt{W_{Kx}^2 + W_{Ky}^2}. \quad (24)$$

The distance from the point K to the instantaneous center of acceleration Q is calculated by the following formula:

$$|KQ| = \frac{W_K}{\sqrt{\omega_2^4 + \varepsilon_2^2}}. \quad (25)$$

To determine the angle between the acceleration W_K and the segment KQ we use the formula.

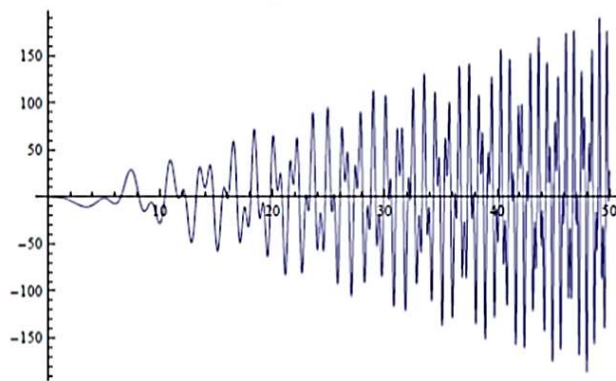
$$tg\theta = \frac{\varepsilon}{\omega^2}. \quad (26)$$

Regarding the given equations and using computer program *Wolfram* numerical examples had been calculated. Based on these examples, changing of the values of speed and acceleration in time in the graphical form are shown below.

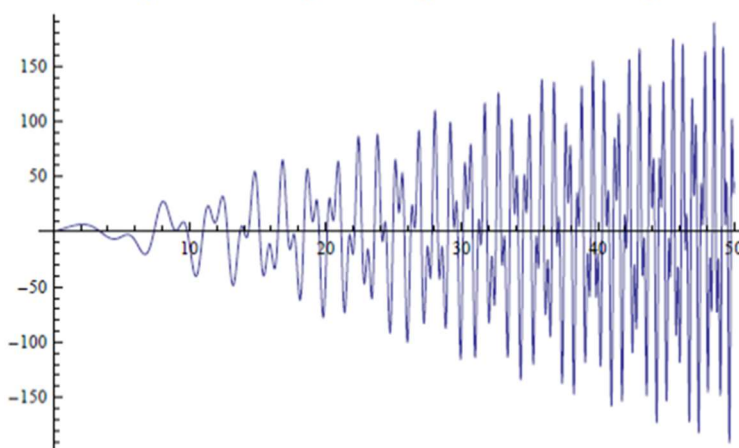
During calculations the angular acceleration ε was taken equal to 0.4 rad/sec; The angular velocity $\omega=0$ when $t=0$. The step during calculations had been taken equal to 0.05 and the time was varied from 0 up to 50 sec.

```
V1[t_] := Plot[-1.6 t x Sin[0.2 x t^2] - 2.311 t x Sin[0.088 x t^2], {t, 0, 50}]
V2[t_] := Plot[1.6 t x Cos[0.2 x t^2] + 2.311 t x Cos[0.088 x t^2], {t, 0, 50}]
V[t_] := Plot[Sqrt[V1^2 + V2^2], {t, 0, 50}, Axes -> {True, True}]
Show[V1, V2]
```

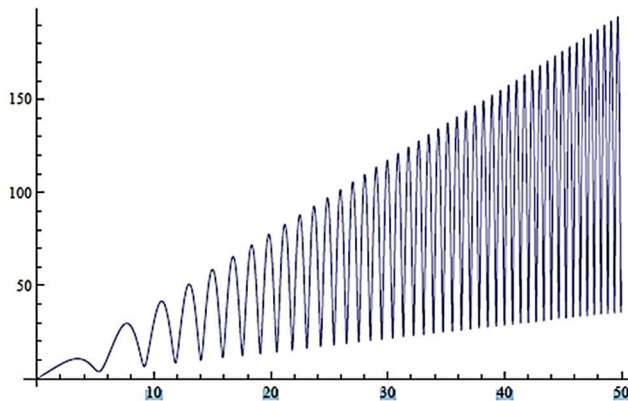
```
V1 = Plot[-1.6 t x Sin[0.2 x t^2] - 2.311 t x Sin[0.088 x t^2], {t, 0, 50}]
```



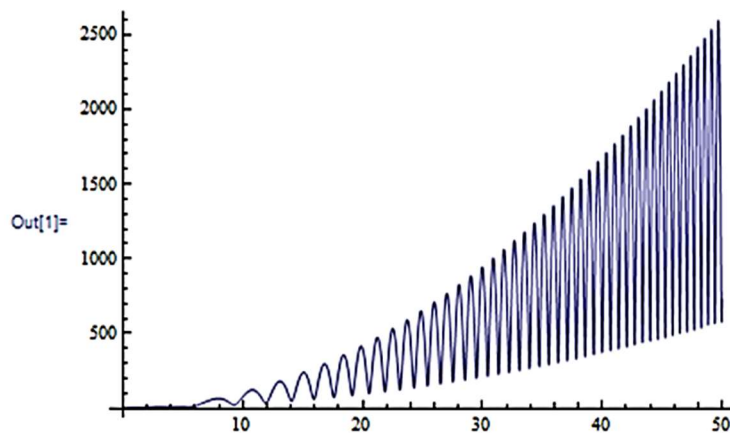
```
V2 = Plot[1.6 t x Cos[0.2 x t^2] + 2.311 t x Cos[0.088 x t^2], {t, 0, 50}]
```



By summing these curves of velocities, we obtain the following summarized plot:

$$V = \text{Plot}\left[\text{Sqrt}\left[\left(1.6 t \times \text{Cos}\left[0.2 \times t^2\right] + 2.311 t \times \text{Cos}\left[0.088 \times t^2\right]\right)^2 + \left(-1.6 t \times \text{Sin}\left[0.2 \times t^2\right] - 2.311 t \times \text{Sin}\left[0.088 \times t^2\right]\right)^2\right], \{t, 0, 50\}\right]$$


In[1]:= W =

$$\text{Plot}\left[\text{Sqrt}\left[\left(-0.64 t^2 \times \text{Cos}\left[0.2 \times t^2\right] - 1.6 t \times \text{Sin}\left[0.2 \times t^2\right] - 0.41 t^2 \times \text{Cos}\left[0.088 \times t^2\right] - 2.311 \text{Sin}\left[0.088 \times t^2\right]\right)^2 + \left(-0.64 t^2 \times \text{Sin}\left[0.2 \times t^2\right] + 1.6 t \times \text{Cos}\left[0.2 \times t^2\right] - 0.41 t^2 \times \text{Sin}\left[0.088 \times t^2\right] + 2.311 \text{Cos}\left[0.088 \times t^2\right]\right)^2\right], \{t, 0, 50\}\right]$$


Conclusion

A kinematic solution of a planetary mechanism with a satellite is presented, which is in external engagement with a fixed central wheel. The geometric features of the mechanism path allow the use of this mechanism in machines, providing technological processes for which

minor oscillations of the driven link are sometimes required at the stop stage. Mechanisms of this type are used in the textile industry, where in the treatment of a material, it is required that the executive body can make slight fluctuations relative to the temporary position.

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UDC 517

SCOPUS CODE 2201

<https://doi.org/10.36073/1512-0996-2024-1-170-175>

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ანოტაცია. სტატიაში წარმოდგენილია სატელიტური პლანეტარული მექანიზმის კინემატიკური კვლევა, რომელიც იმყოფება შიდა მოდელში უძრავ ცენტრალურ კბილა თვალთან. მექანიზმის ტრაექტორიის თავისებურებები საშუალებას იძლევა გამოყენებული იქნეს ეს მექანიზმი მანქანებში, რომლებიც უზრუნველყოფენ ტექნოლოგიურ პროცესებს, რომელთათვისაც საჭიროა მიმყოლი რგოლის მცირე რხევები გაჩერების ეტაპებზე. ასეთი ტიპის მექანიზმები გამოიყენება ტექსტილურ წარმოებაში, სადაც ქსოვილის დამუშავებისას მოითხოვება შემსრულებელი ორგანოს მცირე რხევების დაშვება მექანიზმის დროებითი გაჩერების დროს.

საკვანძო სიტყვები: აჩქარება; მექანიზმი; სატელიტი; სიჩქარე; ტრაექტორია; ციკლოიდა.

The date of review 22.09.2023

The date of submission 24.10.2023

Signed for publishing 22.03.2024